Modified teleparallel theories of gravity f(T,B) Gauss-Bonnet and trace extensions $f(T,B,T_G,B_G,\mathcal{T})$ Conclusions

Modified teleparallel theories of gravity

Sebastián Bahamonde

sebastian.beltran.14@ucl.ac.uk Department of Mathematics, University College London

Quantum Structure of Spacetime and Gravity, Belgrade, Serbia. Work in collaboration with Christian Böhmer 26th August 2016



Modified teleparallel theories of gravity f(T,B) Gauss-Bonnet and trace extensions $f(T,B,T_G,B_G,\mathcal{T})$ Conclusions

Outline



Introduction

- Teleparallel gravity
- f(R) and f(T) gravity
- 2 Modified teleparallel theories of gravity f(T, B)

3 Gauss-Bonnet and trace extensions $f(T, B, T_G, B_G, T)$

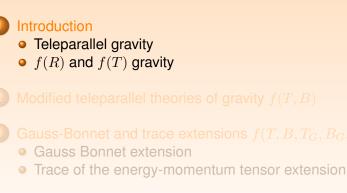
- Gauss Bonnet extension
- Trace of the energy-momentum tensor extension

Onclusions



Modified teleparallel theories of gravity f(T, B)Gauss-Bonnet and trace extensions $f(T, B, T_G, B_G, T)$ Conclusions

Outline



Conclusions

Modified teleparallel theories of gravity f(T, B)Gauss-Bonnet and trace extensions $f(T, B, T_G, B_G, T)$ Conclusions Teleparallel gravity f(R) and f(T) gravity

Teleparallel equivalent of general relativity

• A general connection which defines the parallel transportation is

Most general spin connection (Einstein-Cartan theory)

General spin connection related with curvature related with torsion $\overbrace{w_{\mu}}^{\lambda}{}_{\nu} = \overbrace{\Gamma_{\mu\nu}}^{\lambda} + \overbrace{\tilde{K}_{\mu}}^{\lambda}{}_{\nu}.$

- G.R. assumes for simplicity that $\tilde{K}_{\mu}{}^{\lambda}{}_{\nu} \equiv 0$
- Weitzenböck noticed that it is always possible to define a connection $W_{\mu}{}^{\lambda}{}_{\nu}$ on a space such that is globally flat $(R^{\lambda}{}_{\mu\nu\sigma} \equiv 0) \Longrightarrow \mathsf{TEGR}$

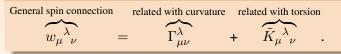


Modified teleparallel theories of gravity f(T, B)Gauss-Bonnet and trace extensions $f(T, B, T_G, B_G, T)$ Conclusions Teleparallel gravity f(R) and f(T) gravity

Teleparallel equivalent of general relativity

• A general connection which defines the parallel transportation is

Most general spin connection (Einstein-Cartan theory)



- G.R. assumes for simplicity that $\tilde{K}_{\mu}{}^{\lambda}{}_{\nu} \equiv 0$
- Weitzenböck noticed that it is always possible to define a connection $W_{\mu}{}^{\lambda}{}_{\nu}$ on a space such that is globally flat $(R^{\lambda}{}_{\mu\nu\sigma} \equiv 0) \Longrightarrow \mathsf{TEGR}$

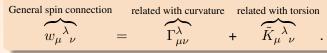


Modified teleparallel theories of gravity f(T, B)Gauss-Bonnet and trace extensions $f(T, B, T_G, B_G, T)$ Conclusions Teleparallel gravity f(R) and f(T) gravity

Teleparallel equivalent of general relativity

• A general connection which defines the parallel transportation is

Most general spin connection (Einstein-Cartan theory)



- G.R. assumes for simplicity that $\tilde{K}_{\mu}{}^{\lambda}{}_{\nu} \equiv 0$
- Weitzenböck noticed that it is always possible to define a connection $W_{\mu}{}^{\lambda}{}_{\nu}$ on a space such that is globally flat $(R^{\lambda}{}_{\mu\nu\sigma} \equiv 0) \Longrightarrow \mathsf{TEGR}$

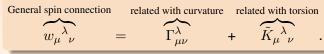


Modified teleparallel theories of gravity f(T, B)Gauss-Bonnet and trace extensions $f(T, B, T_G, B_G, T)$ Conclusions Teleparallel gravity f(R) and f(T) gravity

Teleparallel equivalent of general relativity

• A general connection which defines the parallel transportation is

Most general spin connection (Einstein-Cartan theory)



- G.R. assumes for simplicity that $\tilde{K}_{\mu}{}^{\lambda}{}_{\nu} \equiv 0$
- Weitzenböck noticed that it is always possible to define a connection $W_{\mu}{}^{\lambda}{}_{\nu}$ on a space such that is globally flat $(R^{\lambda}{}_{\mu\nu\sigma} \equiv 0) \Longrightarrow \mathsf{TEGR}$



Modified teleparallel theories of gravity f(T, B)Gauss-Bonnet and trace extensions $f(T, B, T_G, B_G, T)$ Conclusions Teleparallel gravity f(R) and f(T) gravity

Teleparallel equivalent of general relativity

• The teleparallel action is formulated based on a gravitational scalar called the torsion scalar *T*

Here, $e = \det(e_a^{\mu}) = \sqrt{-g}$ and $\kappa = 8\pi G$.

 The relationship between the scalar curvature R and the scalar torsion is

 $R = -T + -\partial_{1}(eT') = -T + B.$

Modified teleparallel theories of gravity f(T, B)Gauss-Bonnet and trace extensions $f(T, B, T_G, B_G, T)$ Conclusions Teleparallel gravity f(R) and f(T) gravity

Teleparallel equivalent of general relativity

• The teleparallel action is formulated based on a gravitational scalar called the torsion scalar *T*

Here, $e = \det(e_a^{\mu}) = \sqrt{-g}$ and $\kappa = 8\pi G$.

 The relationship between the scalar curvature R and the scalar torsion is

 $R = -T + -\partial_{1}(eT') = -T + B.$

Modified teleparallel theories of gravity f(T, B)Gauss-Bonnet and trace extensions $f(T, B, T_G, B_G, T)$ Conclusions Teleparallel gravity f(R) and f(T) gravity

Teleparallel equivalent of general relativity

• The teleparallel action is formulated based on a gravitational scalar called the torsion scalar *T*

TEGR action

$$S_{\text{TEGR}} = \int \left[-\frac{T}{2\kappa} + L_{\text{m}} \right] e \, d^4 x \, .$$

Here,
$$e = \det(e_a^{\mu}) = \sqrt{-g}$$
 and $\kappa = 8\pi G$.

• The relationship between the scalar curvature *R* and the scalar torsion is

Relationship between R and T $R=-T+rac{2}{e}\partial_{\mu}(eT^{\mu})=-T+B.$

Modified teleparallel theories of gravity f(T, B)Gauss-Bonnet and trace extensions $f(T, B, T_G, B_G, T)$ Conclusions Teleparallel gravity f(R) and f(T) gravity

Teleparallel equivalent of general relativity

• The teleparallel action is formulated based on a gravitational scalar called the torsion scalar *T*

TEGR action

$$S_{\text{TEGR}} = \int \left[-\frac{T}{2\kappa} + L_{\text{m}} \right] e \, d^4 x \, .$$

Here, $e = \det(e_a^{\mu}) = \sqrt{-g}$ and $\kappa = 8\pi G$.

• The relationship between the scalar curvature *R* and the scalar torsion is

Relationship between R and I

$$R = -T + \frac{2}{e}\partial_{\mu}(eT^{\mu}) = -T + B.$$

Modified teleparallel theories of gravity f(T, B)Gauss-Bonnet and trace extensions $f(T, B, T_G, B_G, T)$ Conclusions Teleparallel gravity f(R) and f(T) gravity

Teleparallel equivalent of general relativity

• The teleparallel action is formulated based on a gravitational scalar called the torsion scalar *T*

TEGR action

$$S_{\text{TEGR}} = \int \left[-\frac{T}{2\kappa} + L_{\text{m}} \right] e \, d^4 x \, .$$

Here, $e = \det(e_a^{\mu}) = \sqrt{-g}$ and $\kappa = 8\pi G$.

• The relationship between the scalar curvature *R* and the scalar torsion is

Relationship between R and T

$$R = -T + \frac{2}{e}\partial_{\mu}(eT^{\mu}) = -T + B.$$

Modified teleparallel theories of gravity f(T, B)Gauss-Bonnet and trace extensions $f(T, B, T_G, B_G, T)$ Conclusions



Teleparallel gravity f(R) and f(T) gravity

• A well studied modification of GR is *f*(*R*) gravity, which has the following action

$$S_{f(R)} = \int f(R) \sqrt{-g} \, d^4x \, .$$

- Here, *f* is an arbitrary (sufficiently smooth) function of the Ricci scalar.
- Ricci scalar depends on second derivatives of the metric tensor → Fourth order theory.



Modified teleparallel theories of gravity f(T, B)Gauss-Bonnet and trace extensions $f(T, B, T_G, B_G, T)$ Conclusions



Teleparallel gravity f(R) and f(T) gravity

• A well studied modification of GR is *f*(*R*) gravity, which has the following action

$$S_{f(R)} = \int f(R) \sqrt{-g} \, d^4x \, .$$

- Here, *f* is an arbitrary (sufficiently smooth) function of the Ricci scalar.
- Ricci scalar depends on second derivatives of the metric tensor → Fourth order theory.



Modified teleparallel theories of gravity f(T, B)Gauss-Bonnet and trace extensions $f(T, B, T_G, B_G, T)$ Conclusions



Teleparallel gravity f(R) and f(T) gravity

• A well studied modification of GR is *f*(*R*) gravity, which has the following action

$$S_{f(R)} = \int f(R) \sqrt{-g} \, d^4x \, .$$

- Here, *f* is an arbitrary (sufficiently smooth) function of the Ricci scalar.
- Ricci scalar depends on second derivatives of the metric tensor → Fourth order theory.



Modified teleparallel theories of gravity f(T, B)Gauss-Bonnet and trace extensions $f(T, B, T_G, B_G, T)$ Conclusions



Teleparallel gravity f(R) and f(T) gravity

• A well studied modification of GR is *f*(*R*) gravity, which has the following action

$$S_{f(R)} = \int f(R) \sqrt{-g} \, d^4x \, .$$

- Here, *f* is an arbitrary (sufficiently smooth) function of the Ricci scalar.
- Ricci scalar depends on second derivatives of the metric tensor → Fourth order theory.



Modified teleparallel theories of gravity f(T, B)Gauss-Bonnet and trace extensions $f(T, B, T_G, B_G, T)$ Conclusions Teleparallel gravity f(R) and f(T) gravity

f(T) gravity

• In analogy with f(R) gravity, one can consider in the Teleparallel framework, the f(T) gravity

f(T) gravity action

$$S_{f(T)} = \int f(T)e \, d^4x \, .$$

- The torsion scalar T depends on the first derivatives of the tetrads → Second order theory.
- T is not invariant under local LT ⇒ f(T) is also not invariant under local LT.

Modified teleparallel theories of gravity f(T, B)Gauss-Bonnet and trace extensions $f(T, B, T_G, B_G, T)$ Conclusions Teleparallel gravity f(R) and f(T) gravity

f(T) gravity

• In analogy with f(R) gravity, one can consider in the Teleparallel framework, the f(T) gravity

f(T) gravity action

$$S_{f(T)} = \int f(T)e \, d^4x \, .$$

- The torsion scalar T depends on the first derivatives of the tetrads → Second order theory.
- T is not invariant under local LT $\implies f(T)$ is also not invariant under local LT.

Modified teleparallel theories of gravity f(T, B)Gauss-Bonnet and trace extensions $f(T, B, T_G, B_G, T)$ Conclusions Teleparallel gravity f(R) and f(T) gravity

f(T) gravity

• In analogy with f(R) gravity, one can consider in the Teleparallel framework, the f(T) gravity

f(T) gravity action

$$S_{f(T)} = \int f(T)e \, d^4x \, .$$

- The torsion scalar T depends on the first derivatives of the tetrads → Second order theory.
- T is not invariant under local LT ⇒ f(T) is also not invariant under local LT.

Modified teleparallel theories of gravity f(T, B)Gauss-Bonnet and trace extensions $f(T, B, T_G, B_G, T)$ Conclusions Teleparallel gravity f(R) and f(T) gravity

f(T) gravity

• In analogy with f(R) gravity, one can consider in the Teleparallel framework, the f(T) gravity

f(T) gravity action

$$S_{f(T)} = \int f(T)e \, d^4x \, .$$

- The torsion scalar T depends on the first derivatives of the tetrads → Second order theory.
- T is not invariant under local LT ⇒ f(T) is also not invariant under local LT.

Not equivalency between f(R) and f(T)

Modified teleparallel theories of gravity f(T, B)Gauss-Bonnet and trace extensions $f(T, B, T_G, B_G, T)$ Conclusions Teleparallel gravity f(R) and f(T) gravity

f(T) gravity

• In analogy with f(R) gravity, one can consider in the Teleparallel framework, the f(T) gravity

f(T) gravity action

$$S_{f(T)} = \int f(T)e \, d^4x \, .$$

- The torsion scalar T depends on the first derivatives of the tetrads → Second order theory.
- T is not invariant under local LT ⇒ f(T) is also not invariant under local LT.

Not equivalency between f(R) and f(T)

Outline



• f(R) and f(T) gravity

2 Modified teleparallel theories of gravity f(T, B)

- Gauss-Bonnet and trace extensions f(T, B, T_G, B_G, T)
 Gauss Bonnet extension
 - Trace of the energy-momentum tensor extension

Conclusions



f(T,B) gravity

Inspired by the later discussion, we define the action (Bahamonde et. al Phys.Rev. D92,10, 104042 (2015))

f(T,B) gravity action

$$S_{f(T,B)} = \int \left[\frac{1}{\kappa}f(T,B) + L_{\rm m}\right] e \, d^4x \,,$$

where f is a function of both of its arguments and $L_{\rm m}$ is a matter Lagrangian.

• Since R = -T + B, by setting f = f(-T + B) = f(R) we recover f(R) gravity.



f(T,B) gravity

Inspired by the later discussion, we define the action (Bahamonde et. al Phys.Rev. D92,10, 104042 (2015))

f(T,B) gravity action

$$S_{f(T,B)} = \int \left[\frac{1}{\kappa}f(T,B) + L_{\rm m}\right] e \, d^4x \,,$$

where f is a function of both of its arguments and $L_{\rm m}$ is a matter Lagrangian.

• Since R = -T + B, by setting f = f(-T + B) = f(R) we recover f(R) gravity.



f(T,B) gravity

Inspired by the later discussion, we define the action (Bahamonde et. al Phys.Rev. D92,10, 104042 (2015))

f(T,B) gravity action

$$S_{f(T,B)} = \int \left[\frac{1}{\kappa}f(T,B) + L_{\rm m}\right] e \, d^4x \,,$$

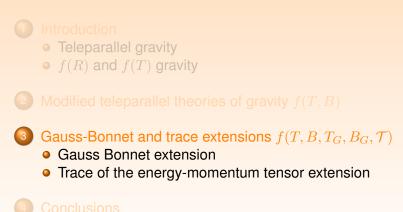
where f is a function of both of its arguments and $L_{\rm m}$ is a matter Lagrangian.

• Since R = -T + B, by setting f = f(-T + B) = f(R) we recover f(R) gravity.



Gauss Bonnet extension Frace of the energy-momentum tensor extension

Outline





10/17

Gauss Bonnet extension Frace of the energy-momentum tensor extension

Gauss-Bonnet extension

• The Gauss-Bonnet term is a quadratic combination of the Riemann tensor and its contractions given by

Gauss-Bonnet term

$$G = R^2 - 4R_{\mu\nu}R^{\mu\nu} + R_{\mu\nu\kappa\lambda}R^{\mu\nu\kappa\lambda}$$

 This term can be expressed in a fashion similar form as before which simply reads

Relationship between Gauss-Bonnet G and T_G

 $G = -T_G + B_G.$



11/17

Gauss Bonnet extension Frace of the energy-momentum tensor extension

Gauss-Bonnet extension

• The Gauss-Bonnet term is a quadratic combination of the Riemann tensor and its contractions given by

Gauss-Bonnet term

$$G = R^2 - 4R_{\mu\nu}R^{\mu\nu} + R_{\mu\nu\kappa\lambda}R^{\mu\nu\kappa\lambda}$$

 This term can be expressed in a fashion similar form as before which simply reads





Gauss Bonnet extension Frace of the energy-momentum tensor extension

Gauss-Bonnet extension

• The Gauss-Bonnet term is a quadratic combination of the Riemann tensor and its contractions given by

Gauss-Bonnet term

$$G = R^2 - 4R_{\mu\nu}R^{\mu\nu} + R_{\mu\nu\kappa\lambda}R^{\mu\nu\kappa\lambda}$$

 This term can be expressed in a fashion similar form as before which simply reads

Relationship between Gauss-Bonnet G and T_G

 $G = -T_G + B_G.$



11/17

Gauss Bonnet extension Frace of the energy-momentum tensor extension

Gauss-Bonnet extension

• The Gauss-Bonnet term is a quadratic combination of the Riemann tensor and its contractions given by

Gauss-Bonnet term

$$G = R^2 - 4R_{\mu\nu}R^{\mu\nu} + R_{\mu\nu\kappa\lambda}R^{\mu\nu\kappa\lambda}$$

 This term can be expressed in a fashion similar form as before which simply reads

Relationship between Gauss-Bonnet G and T_G

$$G = -T_G + B_G.$$



Gauss Bonnet extension Frace of the energy-momentum tensor extension

Gauss-Bonnet extension action

 Inspired by the later discussion, we define the action (Bahamonde et. al arXiv:1606.05557 [gr-qc] (2016))

Gauss-Bonnet extension action

$$S_{f(T,B,T_G,B_G)} = \int \left[\frac{1}{2\kappa}f(T,B,T_G,B_G) + L_{\rm m}\right] e \, d^4x \,.$$

• This action is very general. For instance, we can recover $f = f(T, T_G)$ gravity or $f(-T + B, -T_G + B_G) = f(R, G)$ gravity.



Gauss Bonnet extension Frace of the energy-momentum tensor extension

Gauss-Bonnet extension action

Inspired by the later discussion, we define the action

(Bahamonde et. al arXiv:1606.05557 [gr-qc] (2016))

Gauss-Bonnet extension action

$$S_{f(T,B,T_G,B_G)} = \int \left[\frac{1}{2\kappa}f(T,B,T_G,B_G) + L_{\rm m}\right] e \, d^4x \,.$$

• This action is very general. For instance, we can recover $f = f(T, T_G)$ gravity or $f(-T + B, -T_G + B_G) = f(R, G)$ gravity.



Gauss Bonnet extension Trace of the energy-momentum tensor extension

Gauss-Bonnet extension action

Inspired by the later discussion, we define the action

(Bahamonde et. al arXiv:1606.05557 [gr-qc] (2016))

Gauss-Bonnet extension action

$$S_{f(T,B,T_G,B_G)} = \int \left[\frac{1}{2\kappa}f(T,B,T_G,B_G) + L_{\rm m}\right] e \, d^4x \,.$$

• This action is very general. For instance, we can recover $f = f(T, T_G)$ gravity or $f(-T + B, -T_G + B_G) = f(R, G)$ gravity.



Gauss Bonnet extension Trace of the energy-momentum tensor extension

Trace extension

We will now consider the later framework and include the trace of the energy-momentum tensor to the action discussed before. This gives the extended action

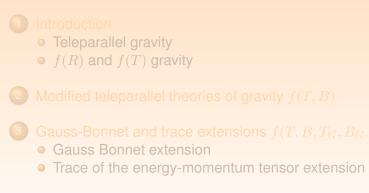
Trace and Gauss-Bonnet action extension

$$S_{f(T,B,T_G,B_G,\mathcal{T})} = \int \left[\frac{1}{2\kappa}f(T,B,T_G,B_G,\mathcal{T}) + L_{\rm m}\right] e \, d^4x \,,$$

where additionally f is a function of the trace of the energy-momentum tensor $\mathcal{T}=E_a^\beta\mathcal{T}_\beta^a$



Outline

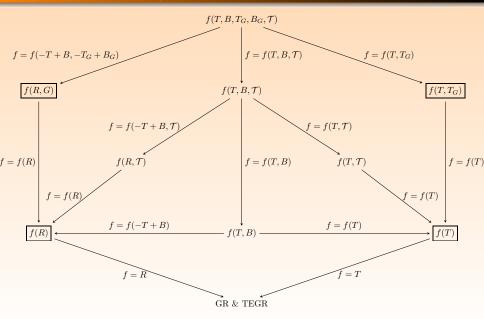






14/17

$f(T, B, T_G, B_G, \mathcal{T})$ diagram



Modified teleparallel theories of gravity f(T, B)Gauss-Bonnet and trace extensions $f(T, B, T_G, B_G, T)$ Conclusions

Conclusions

- For many years now, an ever increasing number of modifications of GR has been considered. Many of these theories were considered in isolation in the past and their relationship with other similarly looking theories was only made implicitly.
- In our works, we explicitly found the corresponding teleparallel equivalent to well-known theories of gravity such as f(R), f(R,G), f(R,T) and others.



Modified teleparallel theories of gravity f(T, B)Gauss-Bonnet and trace extensions $f(T, B, T_G, B_G, T)$ Conclusions

Conclusions

- For many years now, an ever increasing number of modifications of GR has been considered. Many of these theories were considered in isolation in the past and their relationship with other similarly looking theories was only made implicitly.
- In our works, we explicitly found the corresponding teleparallel equivalent to well-known theories of gravity such as f(R), f(R,G), f(R,T) and others.



Modified teleparallel theories of gravity f(T, B)Gauss-Bonnet and trace extensions $f(T, B, T_G, B_G, T)$ Conclusions



- S. Bahamonde, C. G. Böhmer and M. Wright, "Modified teleparallel theories of gravity," Phys. Rev. D **92** (2015) 10, 104042
- S. Bahamonde and C. G. Boehmer, "Modified teleparallel theories of gravity – Gauss-Bonnet and trace extensions," arXiv:1606.05557 [gr-qc].



Modified teleparallel theories of gravity f(T, B)Gauss-Bonnet and trace extensions $f(T, B, T_G, B_G, T)$ Conclusions



- S. Bahamonde, C. G. Böhmer and M. Wright, "Modified teleparallel theories of gravity," Phys. Rev. D **92** (2015) 10, 104042
- S. Bahamonde and C. G. Boehmer, "Modified teleparallel theories of gravity – Gauss-Bonnet and trace extensions," arXiv:1606.05557 [gr-qc].

