[Modified teleparallel theories of gravity](#page-21-0) $f(T, B)$ [Gauss-Bonnet and trace extensions](#page-25-0) $f(T, B, T_G, B_G, \mathcal{T})$ [Conclusions](#page-34-0)

Modified teleparallel theories of gravity

Sebastián Bahamonde

sebastian.beltran.14@ucl.ac.uk Department of Mathematics, University College London

Quantum Structure of Spacetime and Gravity, Belgrade, Serbia. Work in collaboration with Christian Böhmer 26th August 2016

[Modified teleparallel theories of gravity](#page-21-0) $f(T, B)$ [Gauss-Bonnet and trace extensions](#page-25-0) $f(T, B, T_G, B_G, \mathcal{T})$ **[Conclusions](#page-34-0)**

Outline

[Introduction](#page-2-0)

- [Teleparallel gravity](#page-3-0)
- \bullet $f(R)$ and $f(T)$ [gravity](#page-12-0)

[Modified teleparallel theories of gravity](#page-21-0) $f(T, B)$

3 [Gauss-Bonnet and trace extensions](#page-25-0) $f(T, B, T_G, B_G, T)$

- **[Gauss Bonnet extension](#page-26-0)**
- **•** [Trace of the energy-momentum tensor extension](#page-33-0)

[Modified teleparallel theories of gravity](#page-21-0) $f(T, B)$ [Gauss-Bonnet and trace extensions](#page-25-0) $f(T, B, T_G, B_G, T)$ [Conclusions](#page-34-0)

Outline

[Introduction](#page-2-0)

- [Teleparallel gravity](#page-3-0)
- \bullet $f(R)$ and $f(T)$ [gravity](#page-12-0)

• [Gauss Bonnet extension](#page-26-0) • [Trace of the energy-momentum tensor extension](#page-33-0)

[Modified teleparallel theories of gravity](#page-21-0) $f(T, B)$ [Gauss-Bonnet and trace extensions](#page-25-0) $f(T, B, T_G, B_G, \mathcal{T})$ **[Conclusions](#page-34-0)** [Teleparallel gravity](#page-6-0)

Teleparallel equivalent of general relativity

A general connection which defines the parallel transportation is

 $\widetilde{w_{\mu}}^{\lambda}_{\nu} \quad \quad = \quad \quad \widehat{\Gamma}^{\lambda}_{\mu\nu} \qquad \quad + \quad \quad \widetilde{K}_{\mu}{}^{\lambda}_{\nu} \quad \quad .$

- G.R. assumes for simplicity that $\tilde{K}_{\mu}{}^{\lambda}{}_{\nu} \equiv 0$
- Weitzenböck noticed that it is always possible to define a connection $W_\mu{}^\lambda{}_\nu$ on a space such that is globally flat $(R^{\lambda}{}_{\mu\nu\sigma}\equiv 0) \Longrightarrow \mathsf{TEGR}$

[Modified teleparallel theories of gravity](#page-21-0) $f(T, B)$ [Gauss-Bonnet and trace extensions](#page-25-0) $f(T, B, T_G, B_G, \mathcal{T})$ **[Conclusions](#page-34-0)** [Teleparallel gravity](#page-6-0)

Teleparallel equivalent of general relativity

A general connection which defines the parallel transportation is

Most general spin connection (Einstein-Cartan theory)

- G.R. assumes for simplicity that $\tilde{K}_{\mu}{}^{\lambda}{}_{\nu} \equiv 0$
- Weitzenböck noticed that it is always possible to define a connection $W_\mu{}^\lambda{}_\nu$ on a space such that is globally flat $(R^{\lambda}{}_{\mu\nu\sigma}\equiv 0) \Longrightarrow \mathsf{TEGR}$

[Modified teleparallel theories of gravity](#page-21-0) $f(T, B)$ [Gauss-Bonnet and trace extensions](#page-25-0) $f(T, B, T_G, B_G, \mathcal{T})$ **[Conclusions](#page-34-0)** [Teleparallel gravity](#page-6-0)

Teleparallel equivalent of general relativity

A general connection which defines the parallel transportation is

Most general spin connection (Einstein-Cartan theory)

- G.R. assumes for simplicity that $\tilde{K}_{\mu}{}^{\lambda}{}_{\nu} \equiv 0$
- Weitzenböck noticed that it is always possible to define a connection $W_\mu{}^\lambda{}_\nu$ on a space such that is globally flat $(R^{\lambda}{}_{\mu\nu\sigma}\equiv 0) \Longrightarrow \mathsf{TEGR}$

[Modified teleparallel theories of gravity](#page-21-0) $f(T, B)$ [Gauss-Bonnet and trace extensions](#page-25-0) $f(T, B, T_G, B_G, \mathcal{T})$ **[Conclusions](#page-34-0)** [Teleparallel gravity](#page-3-0)

Teleparallel equivalent of general relativity

A general connection which defines the parallel transportation is

Most general spin connection (Einstein-Cartan theory)

- G.R. assumes for simplicity that $\tilde{K}_{\mu}{}^{\lambda}{}_{\nu} \equiv 0$
- Weitzenbock noticed that it is always possible to define a connection $W_\mu{}^\lambda{}_\nu$ on a space such that is globally flat $(R^{\lambda}{}_{\mu\nu\sigma}\equiv 0) \Longrightarrow$ TEGR

[Modified teleparallel theories of gravity](#page-21-0) $f(T, B)$ [Gauss-Bonnet and trace extensions](#page-25-0) $f(T, B, T_G, B_G, T)$ [Conclusions](#page-34-0) [Teleparallel gravity](#page-3-0)

Teleparallel equivalent of general relativity

The teleparallel action is formulated based on a gravitational scalar called the torsion scalar T

Here, $e = \det(e_a^\mu) = \sqrt{-g}$ and $\kappa = 8\pi G$.

[Modified teleparallel theories of gravity](#page-21-0) $f(T, B)$ [Gauss-Bonnet and trace extensions](#page-25-0) $f(T, B, T_G, B_G, T)$ [Conclusions](#page-34-0) [Teleparallel gravity](#page-3-0)

Teleparallel equivalent of general relativity

The teleparallel action is formulated based on a gravitational scalar called the torsion scalar T

Here, $e = \det(e_a^\mu) = \sqrt{-g}$ and $\kappa = 8\pi G$.

[Teleparallel gravity](#page-3-0)

Teleparallel equivalent of general relativity

The teleparallel action is formulated based on a gravitational scalar called the torsion scalar T

TEGR action

$$
S_{\rm TEGR} = \int \left[-\frac{T}{2\kappa} + L_{\rm m} \right] e \, d^4x \,.
$$

Here,
$$
e = \det(e_a^{\mu}) = \sqrt{-g}
$$
 and $\kappa = 8\pi G$.

 \bullet The relationship between the scalar curvature R and the scalar torsion is

[Teleparallel gravity](#page-3-0)

Teleparallel equivalent of general relativity

The teleparallel action is formulated based on a gravitational scalar called the torsion scalar T

TEGR action

$$
S_{\rm TEGR} = \int \left[-\frac{T}{2\kappa} + L_{\rm m} \right] e \, d^4x \,.
$$

Here,
$$
e = \det(e_a^{\mu}) = \sqrt{-g}
$$
 and $\kappa = 8\pi G$.

 \bullet The relationship between the scalar curvature R and the scalar torsion is

$$
R=-T+\frac{2}{e}\partial_{\mu}(eT^{\mu})=-T+B.
$$

[Teleparallel gravity](#page-3-0)

Teleparallel equivalent of general relativity

The teleparallel action is formulated based on a gravitational scalar called the torsion scalar T

TEGR action

$$
S_{\rm TEGR} = \int \left[-\frac{T}{2\kappa} + L_{\rm m} \right] e \, d^4x \,.
$$

Here,
$$
e = \det(e_a^{\mu}) = \sqrt{-g}
$$
 and $\kappa = 8\pi G$.

 \bullet The relationship between the scalar curvature R and the scalar torsion is

Relationship between R and T

$$
R=-T+\frac{2}{e}\partial_{\mu}(eT^{\mu})=-T+B.
$$

[Modified teleparallel theories of gravity](#page-21-0) $f(T, B)$ [Gauss-Bonnet and trace extensions](#page-25-0) $f(T, B, T_G, B_G, \mathcal{T})$ **[Conclusions](#page-34-0)**

 $f(R)$ and $f(T)$ [gravity](#page-15-0)

A well studied modification of GR is $f(R)$ gravity, which has the following action

$$
S_{f(R)} = \int f(R) \sqrt{-g} d^4x.
$$

- \bullet Here, f is an arbitrary (sufficiently smooth) function of the Ricci scalar.
- Ricci scalar depends on second derivatives of the metric tensor → **Fourth order theory**.

[Modified teleparallel theories of gravity](#page-21-0) $f(T, B)$ [Gauss-Bonnet and trace extensions](#page-25-0) $f(T, B, T_G, B_G, \mathcal{T})$ **[Conclusions](#page-34-0)**

 $f(R)$ and $f(T)$ [gravity](#page-15-0)

A well studied modification of GR is $f(R)$ gravity, which has the following action

 $f(R)$ gravity action

$$
S_{f(R)}=\int f(R)\sqrt{-g}\,d^4x\,.
$$

- \bullet Here, f is an arbitrary (sufficiently smooth) function of the Ricci scalar.
- Ricci scalar depends on second derivatives of the metric tensor → **Fourth order theory**.

[Modified teleparallel theories of gravity](#page-21-0) $f(T, B)$ [Gauss-Bonnet and trace extensions](#page-25-0) $f(T, B, T_G, B_G, \mathcal{T})$ **[Conclusions](#page-34-0)**

 $f(R)$ and $f(T)$ [gravity](#page-15-0)

• A well studied modification of GR is $f(R)$ gravity, which has the following action

 $f(R)$ gravity action

$$
S_{f(R)}=\int f(R)\sqrt{-g}\,d^4x\,.
$$

- \bullet Here, f is an arbitrary (sufficiently smooth) function of the Ricci scalar.
- Ricci scalar depends on second derivatives of the metric tensor → **Fourth order theory**.

[Modified teleparallel theories of gravity](#page-21-0) $f(T, B)$ [Gauss-Bonnet and trace extensions](#page-25-0) $f(T, B, T_G, B_G, \mathcal{T})$ [Conclusions](#page-34-0)

 $f(R)$ and $f(T)$ [gravity](#page-12-0)

• A well studied modification of GR is $f(R)$ gravity, which has the following action

 $f(R)$ gravity action

$$
S_{f(R)}=\int f(R)\sqrt{-g}\,d^4x\,.
$$

- \bullet Here, f is an arbitrary (sufficiently smooth) function of the Ricci scalar.
- **Ricci scalar depends on second derivatives of the metric** tensor → **Fourth order theory**.

[Modified teleparallel theories of gravity](#page-21-0) $f(T, B)$ [Gauss-Bonnet and trace extensions](#page-25-0) $f(T, B, T_G, B_G, \mathcal{T})$ **[Conclusions](#page-34-0)** $f(R)$ and $f(T)$ [gravity](#page-12-0)

$f(T)$ gravity

• In analogy with $f(R)$ gravity, one can consider in the Teleparallel framework, the $f(T)$ gravity

$$
S_{f(T)} = \int f(T) e d^4x.
$$

- \bullet The torsion scalar T depends on the first derivatives of the tetrads → **Second order theory**.
- T is not invariant under local $LT \implies f(T)$ is also not **invariant under local LT.**

[Modified teleparallel theories of gravity](#page-21-0) $f(T, B)$ [Gauss-Bonnet and trace extensions](#page-25-0) $f(T, B, T_G, B_G, \mathcal{T})$ **[Conclusions](#page-34-0)** $f(R)$ and $f(T)$ [gravity](#page-12-0)

$f(T)$ gravity

• In analogy with $f(R)$ gravity, one can consider in the Teleparallel framework, the $f(T)$ gravity

 $f(T)$ gravity action

$$
S_{f(T)} = \int f(T)e d^4x.
$$

- \bullet The torsion scalar T depends on the first derivatives of the tetrads → **Second order theory**.
- T is not invariant under local $LT \implies f(T)$ is also not **invariant under local LT.**

[Modified teleparallel theories of gravity](#page-21-0) $f(T, B)$ [Gauss-Bonnet and trace extensions](#page-25-0) $f(T, B, T_G, B_G, \mathcal{T})$ **[Conclusions](#page-34-0)** $f(R)$ and $f(T)$ [gravity](#page-12-0)

$f(T)$ gravity

• In analogy with $f(R)$ gravity, one can consider in the Teleparallel framework, the $f(T)$ gravity

 $f(T)$ gravity action

$$
S_{f(T)} = \int f(T)e d^4x.
$$

- \bullet The torsion scalar T depends on the first derivatives of the tetrads → **Second order theory**.
- T is not invariant under local $LT \Longrightarrow f(T)$ is also not **invariant under local LT.**

[Modified teleparallel theories of gravity](#page-21-0) $f(T, B)$ [Gauss-Bonnet and trace extensions](#page-25-0) $f(T, B, T_G, B_G, \mathcal{T})$ **[Conclusions](#page-34-0)** $f(R)$ and $f(T)$ [gravity](#page-12-0)

$f(T)$ gravity

• In analogy with $f(R)$ gravity, one can consider in the Teleparallel framework, the $f(T)$ gravity

 $f(T)$ gravity action

$$
S_{f(T)} = \int f(T)e d^4x.
$$

- \bullet The torsion scalar T depends on the first derivatives of the tetrads → **Second order theory**.
- T is not invariant under local $LT \implies f(T)$ is also not **invariant under local LT.**

Field equations of $f(T) \neq$ Field equations of $f(R)$

[Modified teleparallel theories of gravity](#page-21-0) $f(T, B)$ [Gauss-Bonnet and trace extensions](#page-25-0) $f(T, B, T_G, B_G, \mathcal{T})$ **[Conclusions](#page-34-0)** $f(R)$ and $f(T)$ [gravity](#page-12-0)

$f(T)$ gravity

• In analogy with $f(R)$ gravity, one can consider in the Teleparallel framework, the $f(T)$ gravity

 $f(T)$ gravity action

$$
S_{f(T)} = \int f(T)e d^4x.
$$

- \bullet The torsion scalar T depends on the first derivatives of the tetrads → **Second order theory**.
- T is not invariant under local $LT \implies f(T)$ is also not **invariant under local LT.**

Not equivalency between $f(R)$ and $f(T)$

Field equations of $f(T) \neq$ Field equations of $f(R)$

Outline

[Modified teleparallel theories of gravity](#page-21-0) $f(T, B)$

• [Gauss Bonnet extension](#page-26-0) • [Trace of the energy-momentum tensor extension](#page-33-0)

$f(T, B)$ gravity

• Inspired by the later discussion, we define the action (Bahamonde et. al Phys.Rev. D92,10, 104042 (2015))

$$
S_{f(T,B)} = \int \left[\frac{1}{\kappa}f(T,B) + L_{\rm m}\right] e d^4x,
$$

where f is a function of both of its arguments and L_m is a matter Lagrangian.

• Since $R = -T + B$, by setting $f = f(-T + B) = f(R)$ we recover $f(R)$ gravity.

$f(T, B)$ gravity

• Inspired by the later discussion, we define the action (Bahamonde et. al Phys.Rev. D92,10, 104042 (2015))

$f(T, B)$ gravity action

$$
S_{f(T,B)} = \int \left[\frac{1}{\kappa}f(T,B) + L_{\rm m}\right] e d^4x,
$$

where f is a function of both of its arguments and L_m is a matter Lagrangian.

• Since $R = -T + B$, by setting $f = f(-T + B) = f(R)$ we recover $f(R)$ gravity.

$f(T, B)$ gravity

• Inspired by the later discussion, we define the action (Bahamonde et. al Phys.Rev. D92,10, 104042 (2015))

$f(T, B)$ gravity action

$$
S_{f(T,B)} = \int \left[\frac{1}{\kappa}f(T,B) + L_{\rm m}\right] e d^4x,
$$

where f is a function of both of its arguments and L_m is a matter Lagrangian.

• Since $R = -T + B$, by setting $f = f(-T + B) = f(R)$ we recover $f(R)$ gravity.

Outline

- **•** [Teleparallel gravity](#page-3-0)
- \bullet $f(R)$ and $f(T)$ [gravity](#page-12-0)
-

3 [Gauss-Bonnet and trace extensions](#page-25-0) $f(T, B, T_G, B_G, T)$

- **[Gauss Bonnet extension](#page-26-0)**
- **•** [Trace of the energy-momentum tensor extension](#page-33-0)

[Gauss Bonnet extension](#page-29-0)

Gauss-Bonnet extension

The Gauss-Bonnet term is a quadratic combination of the Riemann tensor and its contractions given by

$$
G = R^2 - 4R_{\mu\nu}R^{\mu\nu} + R_{\mu\nu\kappa\lambda}R^{\mu\nu\kappa\lambda}
$$

This term can be expressed in a fashion similar form as before which simply reads

$$
G=-T_G+B_G.
$$

[Gauss Bonnet extension](#page-29-0)

.

Gauss-Bonnet extension

The Gauss-Bonnet term is a quadratic combination of the Riemann tensor and its contractions given by

Gauss-Bonnet term

$$
G = R^2 - 4R_{\mu\nu}R^{\mu\nu} + R_{\mu\nu\kappa\lambda}R^{\mu\nu\kappa\lambda}
$$

This term can be expressed in a fashion similar form as before which simply reads

$$
G=-T_G+B_G.
$$

[Gauss Bonnet extension](#page-29-0)

.

Gauss-Bonnet extension

The Gauss-Bonnet term is a quadratic combination of the Riemann tensor and its contractions given by

Gauss-Bonnet term

$$
G = R^2 - 4R_{\mu\nu}R^{\mu\nu} + R_{\mu\nu\kappa\lambda}R^{\mu\nu\kappa\lambda}
$$

This term can be expressed in a fashion similar form as before which simply reads

$$
G=-T_G+B_G.
$$

[Gauss Bonnet extension](#page-26-0)

.

Gauss-Bonnet extension

The Gauss-Bonnet term is a quadratic combination of the Riemann tensor and its contractions given by

Gauss-Bonnet term

$$
G = R^2 - 4R_{\mu\nu}R^{\mu\nu} + R_{\mu\nu\kappa\lambda}R^{\mu\nu\kappa\lambda}
$$

This term can be expressed in a fashion similar form as before which simply reads

Relationship between Gauss-Bonnet G and T_G

$$
G=-T_G+B_G.
$$

[Gauss Bonnet extension](#page-26-0)

Gauss-Bonnet extension action

• Inspired by the later discussion, we define the action (Bahamonde et. al arXiv:1606.05557 [gr-qc] (2016))

$$
S_{f(T,B,T_G,B_G)} = \int \left[\frac{1}{2\kappa}f(T,B,T_G,B_G) + L_m\right]e d^4x.
$$

• This action is very general. For instance, we can recover $f = f(T, T_G)$ gravity or $f(-T + B, -T_G + B_G) = f(R, G)$ gravity.

[Gauss Bonnet extension](#page-26-0)

Gauss-Bonnet extension action

• Inspired by the later discussion, we define the action

(Bahamonde et. al arXiv:1606.05557 [gr-qc] (2016))

Gauss-Bonnet extension action

$$
S_{f(T,B,T_G,B_G)} = \int \left[\frac{1}{2\kappa} f(T,B,T_G,B_G) + L_m \right] e d^4x.
$$

• This action is very general. For instance, we can recover $f = f(T, T_G)$ gravity or $f(-T + B, -T_G + B_G) = f(R, G)$ gravity.

[Gauss Bonnet extension](#page-26-0)

Gauss-Bonnet extension action

• Inspired by the later discussion, we define the action

(Bahamonde et. al arXiv:1606.05557 [gr-qc] (2016))

Gauss-Bonnet extension action

$$
S_{f(T,B,T_G,B_G)} = \int \left[\frac{1}{2\kappa} f(T,B,T_G,B_G) + L_m \right] e d^4x.
$$

• This action is very general. For instance, we can recover $f = f(T, T_G)$ gravity or $f(-T + B, -T_G + B_G) = f(R, G)$ gravity.

[Trace of the energy-momentum tensor extension](#page-33-0)

Trace extension

We will now consider the later framework and include the trace of the energy-momentum tensor to the action discussed before. This gives the extended action

Trace and Gauss-Bonnet action extension

$$
S_{f(T,B,T_G,B_G,\mathcal{T})} = \int \left[\frac{1}{2\kappa}f(T,B,T_G,B_G,\mathcal{T}) + L_m\right] e d^4x,
$$

where additionally f is a function of the trace of the energy-momentum tensor $\mathcal{T}=E^{\beta}_{a} \mathcal{T}^{a}_{\beta}$

Outline

• [Gauss Bonnet extension](#page-26-0) • [Trace of the energy-momentum tensor extension](#page-33-0)

$f(T, B, T_G, B_G, \mathcal{T})$ diagram

[Modified teleparallel theories of gravity](#page-21-0) $f(T, B)$ [Gauss-Bonnet and trace extensions](#page-25-0) $f(T, B, T_G, B_G, \mathcal{T})$ [Conclusions](#page-34-0)

Conclusions

- **For many years now, an ever increasing number of** modifications of GR has been considered. Many of these theories were considered in isolation in the past and their relationship with other similarly looking theories was only made implicitly.
- In our works, we explicitly found the corresponding teleparallel equivalent to well-known theories of gravity such as $f(R)$, $f(R, G)$, $f(R, \mathcal{T})$ and others.

[Modified teleparallel theories of gravity](#page-21-0) $f(T, B)$ [Gauss-Bonnet and trace extensions](#page-25-0) $f(T, B, T_G, B_G, \mathcal{T})$ [Conclusions](#page-34-0)

Conclusions

- **For many years now, an ever increasing number of** modifications of GR has been considered. Many of these theories were considered in isolation in the past and their relationship with other similarly looking theories was only made implicitly.
- In our works, we explicitly found the corresponding teleparallel equivalent to well-known theories of gravity such as $f(R)$, $f(R, G)$, $f(R, \mathcal{T})$ and others.

[Modified teleparallel theories of gravity](#page-21-0) $f(T, B)$ [Gauss-Bonnet and trace extensions](#page-25-0) $f(T, B, T_G, B_G, T)$ **[Conclusions](#page-34-0)**

- S. Bahamonde, C. G. Böhmer and M. Wright, "Modified" H teleparallel theories of gravity," Phys. Rev. D **92** (2015) 10, 104042
- F

[Modified teleparallel theories of gravity](#page-21-0) $f(T, B)$ [Gauss-Bonnet and trace extensions](#page-25-0) $f(T, B, T_G, B_G, \mathcal{T})$ **[Conclusions](#page-34-0)**

Bibliography

- S. Bahamonde, C. G. Böhmer and M. Wright, "Modified" H teleparallel theories of gravity," Phys. Rev. D **92** (2015) 10, 104042
- F. S. Bahamonde and C. G. Boehmer, "Modified teleparallel theories of gravity – Gauss-Bonnet and trace extensions," arXiv:1606.05557 [gr-qc].

